

Start: ex of irreversible proc (study into why & how it; why, what kind of work is done by friction heat etc) not rep'd on equl spec of states as therm proc, as intermediate states are infinitely nonequil or finally can prob to relaxation to equl states

- irreversible proc rep'd on equl spec of states? $\Delta S > 0$ for whole system
 slow friction - heating? No - all curves reversible on equl spec of states $\Delta S = 0$ if $\Delta S_{\text{univ}} > 0$ is isolated
 need a Comp. for system SOS for irreversible proc -

irrevers: heat flow by conduction from A to B; other no trans in which $\Delta S_{\text{univ}} > 0$ w/ some initial cond, no other changes occur anywhere

role of time for curves on equl SOS

- parametric, curves the only thing distinguish 'directions' of curves are signs of ΔQ & ΔW

first, sketch "naive meaning" of curve on equl SOS, draw pic

- isochore
- isothermal
- isentropic
- adiabatic (reversible, no heat exchange)
- cycle - returns to beginning, though not any change

"textbook" view of therm proc

an equl curve on SOS represents a "series of steps" of infinitesimally small changes occurring infinitely slowly - one needs a driving force to push one out of equl, Min. force has to "push one" back to equl \Rightarrow no net reversible proc in nature, but excellent approx to that 'slow' ones, w/ debatably changing parameters, as such - in limit of no small changes, total proc takes so long

Norton

- his diagnosis of problem:

Problem:

- 1) therm rev proc have nonequil initial & driving forces, otherwise no change of state
- 2) they are sets of equl states in end of which - ΔS in balance of forces

usual resolutions fail:

- no slow proc never progresses
 - 'insensibly small' diff from equl is still not equl
- \Rightarrow either sys changes or it doesn't!

Lecture: Fads Thm of SML - Thermo Props (Norton & Solente)

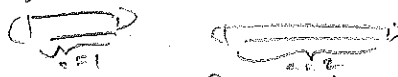
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Norton on ideal vs approx (∞ limit)

- consider 'limit' of systems S_1, S_2, \dots that "approach" (in some sense) limit sys S_∞ , w/ corresponding properties P_1, P_2, \dots that also 'approach' some limit property P_∞
 - 1) it may happen that \exists well defined S_∞ , w/ that P_∞ may still 'exist'
 - 2) even if so, \exists may \sim well defined property P_∞ , or it may not "behave properly"

- if \exists well defined S_∞ , that is idealization = "real or fictitious sys, distinct from target sys" where properties provide inexact desc of some aspects of target sys

sequence of spheres \Rightarrow cylinder



ratio of surf area/vol = $(2\pi a + 4\pi)/(4\pi a^3/3)$

$\Rightarrow 2 \rightarrow a \rightarrow \infty$

agrees with simple limit of ∞ sphere cylinder

- if one takes limit of spheres w/ rad $a=1, 2, 3, \dots$, one gets no limit sys

- \exists may limit sys w/ 'wrong' limit prop
 unit sphere growing into ellipsoids w/ semi-major axis $a = 2, 3, 4, \dots$
 (unit semi-minor axis)



"grows to cylinder", but now surf area/vol ratio is $(\pi^2 a)/(4\pi a/3) = 3\pi/4 \rightarrow a \rightarrow \infty$

\Rightarrow approx = "an inexact desc of target system, properties that approximate those of target system"

- also, approx w/no limit system - e.g., " ∞ box of gas" w/ ratio N/V (particle #/vol) held constant in ∞ limit

→ resolution to paradox: thermodynamics processes are approximations not idealizations

~ inexact propositional descriptions of systems of interest

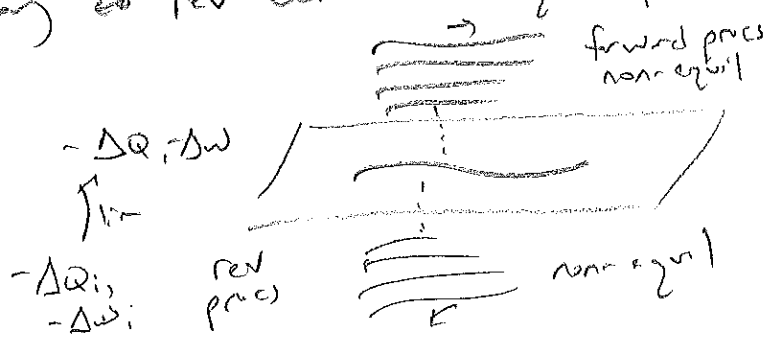
construction

- rep revers proc by ^{continuous} sequence of irrevers proc
- "limiting" set of equil states is rev proc (curve in eqs)
- properties of "rev proc" recovered by limits of props of in sequence of irrev proc
- provides properties only, not a system/proc being them, since \exists such a limiting proc as "real, target proc"
- ~ approx, not an ideal

properties characterizing sequence ~~of limiting states about forward vs reverse process, not about the process~~

- 1) each proc may exchange heat/work w/ env, w/ \exists of F & in behind forces
- 2) proc. with 'forward' or 'reverse' depending on sign of ΔQ and ΔW
- 3) in each sequence, \exists proc in which driving forces are as small as one pleases
- 4) in limit, driving forces $\rightarrow 0$
- 5) limiting values of ΔQ and ΔW are of opposite sign for reverse proc ^{equil. state}
- 6) limiting qs of ΔQ & ΔW at n -th stage of proc = those computed along eq rev curve in equil space

Picture



limit $\Delta Q_i, \Delta W_i$ $\left\{ \begin{array}{l} \text{"limit of heat transfer"} \\ \text{w/o temp diff"} \end{array} \right.$

Valente

- 2 ways out of paradox

- 1) either admit rev proc as pass through non-equil states, but deviations from equil curves are in some sense negligible
- 2) deny any rev proc as equil curve is therm proc

Norton takes 2nd Valente (cf. la Afonso isjawa) takes 2

- for 2 one needs:
- a) show that rev proc can be constructed as equil curve in 'quasi-static' limit
 - b) connect construction to 'real' therm proc

Criticism of Norton

- must be careful: ∞ -time limit of non-equil curve doesn't always yield rev proc (P_{12} has no heat exchange, P_{21} does); depends on how one takes limit
- just coz ∞ -time limit has contradictory properties, it doesn't follow that rev proc does

$\rightarrow P_{21}$ has $\Delta Q = 0$ and $\Delta T = 0$, so it is an ideal in Norton's sense! (no contradictory/impossible properties) it's just not an actual process (cf. no cylinder)

\Rightarrow saying rev proc is "approx" not "ideal" does no real work

- how to define "close enough to equil"? No topology on sequence of non-equil curves, coz no larger space in which they and equil states are embedded, not in classical thermo
- \Rightarrow Giovanni grants Norton too much on bottom of 2nd page of §4.1 - can't define "continuous sequence of curves" and can't integrate along them, coz no \mathbb{R} space, and so no topology and no differential and no affine struc (needed for integration)

- Norton's account leaves out processes that start and end in equilibrium states (as real thermodynamic processes often do), since all non-equilibrium curves in his sequence contain no equilibrium states
→ this seems to me easy to rectify

My criticism of Norton

- whence comes the non-equilibrium curves?
→ imports external machinery from outside classical thermodynamics, no account of how they relate to, fit with, thermodynamic structure
- there may be no ~~such~~ such cases or may not be unique; in latter case, how to decide? (thermodynamics can't tell us!)
- thus, physical significance of limit is not clear
- ~~definition~~ ^{of} ~~propos~~ ^{therms} out of equilibrium not well defined, and contra Norton (p. 49) this is not a problem for 'standard accounts'
- I assumption (set of non-equilibrium curves w/ rev process as limit) holds only in "frictionless" world in which ordinary thermodynamics holds at all scales" (p. 50) ⇒ unsatisfactory for analyzing empirical content of classical thermodynamics

Valente's positive account

- following Afanassjewa, deny rev processes are thermodynamic at all but rather only a "mere mathematical construction useful for calculations" using differential calculus

- Call, new proc 'quasi-proc' in so far as not realistic
- in any real proc - just oriented curve on equil SDS - pure math
- need now to show connection to real thermo proc

⇒ quasi-state proc - a sys can be made to assume any equil state by external forces, and such can be used to "approx" quasi-process by a dense, discrete sequence of equil states (part of quasi-proc) such that during proc connecting any two equil states, state of sys is "close to" equil state

⇒ "iterated equilibria", pass through intermediary non-equil states, an elementary quasi-state proc

- when driving forces go to zero, dense discrete seq of equil states goes to quasi-process (curve on equil SDS)

to make correct w/ real thermo proc

- "∞ time limit" means "sufficiently slowly that in each state we may neglect deviation from equil": "∞ slow ≠ quasi-state"
- ⇒ so quasi-state complete in finite time
- init and final states are equil states

defense:

- "close enough to" equil states? ⇒ 'fills length perceptible errors'

My criticism

- same problem as w/ Norton re: topology, spec, diff struc
- ⇒ quasi-state proc are not in equil SDS, so what does "close enough" mean precisely? how to rep it in machinery of pure classical therm?

Thermodynamical Processes as Equilibrium Processes

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Contents

1. recent controversy: [Norton \(2016\)](#) (wrong, along with Duhem 1903 that Norton cites, for several reasons) versus [Valente \(2017\)](#) (better than Norton, but still faces problems: in what space do the “elementary quasi-static processes” live? what topology defines “closeness” to reversible paths in equilibrium space of states?)
2. I propose an analysis that is consistent with the orthodoxy (*e.g.*, [Fermi \(1956\)](#)) that I think Norton rightly draws attention to as potentially inconsistent, whereas mine has not even the whiff of inconsistency, and has no need for resources external to classical thermodynamics (*e.g.*, non-equilibrium states):
 - a. “equilibrium” is a concept that is always relative to a time scale (*e.g.*, the evaporation of the glass of water—the water is in equilibrium over scales of a few minutes but not several hours), and to an acuity of proposed types of measurement relevant to envisaged experiments
 - b. choose a small enough time scale that the system is essentially in equilibrium over it, say τ units of time
 - c. for the process at issue, determine the equilibrium states at discrete intervals of τ units of time
 - d. this gives a discrete set of points in the equilibrium space of states, labeled by instants of time every τ units
 - e. now, construct an interpolated continuous limit curve between the discrete points, as follows: just as equilibrium is relative to time scale, and so sets lower bounds on the admissible acuity of temporal measurements, so the same holds for the admissible acuity of the measurements of all thermodynamical quantities. For the pragmatic purposes of the proposed experiment, therefore, given two equilibrium states between which all relevant quantities differ by at most the admissible acuities, those two states for all intents and purposes “differ by an infinitesimal”, and so a linear interpolation (“tangent vector” on the space of states) between them is justified.

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- f. label the points of the full curve by the appropriate instants of time (these labels are not part of the curve in the equilibrium space of states *per se*, but an externally imposed device), *viz.*, by a parametrization with respect to the changes of the relevant physical quantities are linear—the integrated time of the complete curve will be finite
- g. one can now integrate quantities (total heat exchanged, total work performed, *etc.*) along the curve in the usual way

One fundamental question for may account: does its pragmatic aspect, in which each curve must be treated differently depending on what experiment one envisages (and so what the time-scale for equilibrium is, and what acuties are admissible in differences of quantities), suffice for a foundational (or “philosophical”) explication? I think so. The only reason it could not is because one holds to the idea that the empirical meaning of theoretical terms and structures in a theory is exhausted by their “ontological” import, irrespective of how they are instantiated in particular experiments. The goal, however, after all, is to tease out, analyze, and characterize the *empirical* content of the theory by an explication of the theoretical structures in a way consonant with their possible experimental applications, and indeed so as to be able to model and explain the possible experimental applications. Such an explication must also be capable of explaining the possibilities of our broader, general “experience” of the kinds of physical systems the theory treats, *e.g.*, to explain why fluids do not maintain shape without external constraint. That one can do this in isolation from consideration of how the theory is applied in practice to actual experiments is a highly non-trivial assumption, though one endemic to contemporary philosophy of science in general, and philosophy of physics in particular—but one which I deplore and reject. (See, *e.g.*, Curiel 2014, 2017a, 2017b.)

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